

Integral representation of the modified Bessel function:  
(slide 24)

$$K_\lambda(z) = \frac{\pi}{4} \int_0^\infty dt \frac{1}{t^2} e^{-\left(t + \frac{z^2}{4t}\right)}$$

Show

$$K_\lambda(z) = \sqrt{\frac{\pi}{2z}} e^{-z} \quad \text{for } z \rightarrow \infty$$

Solution: Perform a saddle point expansion for the integral

$$J(t) = \int_0^\infty dt \frac{1}{t^2} e^{-\left(t + \frac{z^2}{4t}\right)}$$

$$=: \int_0^\infty dt e^{-\bar{F}(t)}$$

$$\bar{F}(t) = \ln t^2 + t + \frac{z^2}{4t}$$

Find extremal value  $\bar{t}$  with  $\bar{F}'(\bar{t})=0$ ,  $\bar{F}''(\bar{t}) > 0$  (local minimum)

$$\bar{F}'(t) = \frac{2}{t} + 1 - \frac{z^2}{4} \frac{1}{t^2}$$

$$\bar{F}'(\bar{t}) = 0 \Rightarrow \bar{t}_\pm = -1 \pm \sqrt{1 + \frac{z^2}{4}}$$

For  $z$  large we find approx.:  $\bar{t} \approx \frac{z}{2}$

$$\bar{F}''(t) = -\frac{2}{t^2} + \frac{z^2}{2} \frac{1}{t^3}$$

$$\bar{F}''(\bar{t}) = -2 \frac{4}{z^2} + \frac{z^2}{2} \cdot \frac{8}{z^3} = -\frac{8}{z^2} + \frac{4}{z} \approx \frac{4}{z} \quad \text{for } z \rightarrow \infty$$

We also need

$$\begin{aligned} \bar{F}(\bar{t}) &= \ln \frac{z^2}{4} + \frac{z}{2} + \frac{z^2}{4} \frac{z}{2} \\ &= \ln \frac{z^2}{4} + \frac{z}{2} + \frac{z}{2} = \ln \frac{z^2}{4} + z \end{aligned}$$

Approximation :

$$\begin{aligned}
 J(t) &= \int_0^\infty dt e^{-F(t)} \\
 &\approx \int_{-\infty}^{+\infty} dt e^{-F(\bar{t}) - \frac{1}{2} F''(\bar{t}) (t - \bar{t})^2} \\
 &= e^{-F(\bar{t})} \underbrace{\int_{-\infty}^{+\infty} dt e^{-\frac{1}{2} F''(\bar{t}) (t - \bar{t})^2}}_{\text{Gauss integral}}
 \end{aligned}$$

a)  $e^{-F(\bar{t})} = e^{-(\ln \frac{z^2}{4} + z)} = \frac{4}{z^2} \cdot e^{-z}$

b)  $\int_{-\infty}^{+\infty} dt e^{-\alpha t^2} = \sqrt{\frac{\pi}{\alpha}} = \sqrt{\frac{\pi}{2} z}$   
 $\alpha = \frac{2}{z}$

$\Rightarrow k_n(z) \propto \frac{z}{4} \cdot \frac{4}{z^2} e^{-z} \cdot \sqrt{\frac{\pi}{2} z} = \sqrt{\frac{\pi}{2z}} e^{-z}$

$O(4)$  breaking at  $O(\alpha^2)$ :  $\mathcal{L}_{\text{break}} = c \cdot \alpha^2 \bar{q} \gamma_\mu D_\mu D_\mu q$

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- breaks  $O(4)$
- preserves chiral symmetry

Map to ChPT:

Promote  $c \cdot \alpha^2$  to a tensor field  $C_{\alpha\beta\gamma\delta}$

which transforms as a singlet under chiral transformations

$$C_{\alpha\beta\gamma\delta} \xrightarrow{R_{\nu L}} C_{\alpha\beta\gamma\delta}$$

"Physical value":  $C_{\alpha\beta\gamma\delta} = \sum_i S_{\alpha\mu} S_{\beta\mu} S_{\gamma\mu} S_{\delta\mu}$

↪

$$\mathcal{L}_{\text{break}} = \bar{q} C_{\alpha\beta\gamma\delta} \gamma_\alpha D_\beta D_\gamma D_\delta q$$

is  $O(4)$  invariant if  $C_{\alpha\beta\gamma\delta}$  is an  $O(4)$  tensor, rank 4

⇒ Construct invariants with  $C_{\alpha\beta\gamma\delta}$  in ChPT!

↪ We need chirally invariant terms with 4  $O(4)$  indices!

↪

$$X_1 = \text{tr } \partial_\alpha U \partial_\beta U^\dagger \partial_\gamma U \partial_\delta U^\dagger$$

$$X_2 = \text{tr } \partial_\alpha U^\dagger \partial_\beta U \partial_\gamma U^\dagger \partial_\delta U$$

Want:  $X_1 \xrightarrow{P} X_2 \Rightarrow$  We need  $X_1 + X_2$  to be  $P$  invariant

There are no additional independent terms:

e.g.

$$\text{tr } \partial_\alpha U \partial_\beta U^+ \underbrace{U \partial_\gamma U^+}_{\text{LEC}} \underbrace{U \partial_\delta U^+}_{\text{LEC}} = -x_1 \\ = -\partial_\gamma U U^+$$

$$\text{tr } \partial_\alpha U \underbrace{U^+ \partial_\beta U}_{\text{LEC}} \underbrace{\partial_\gamma U^+}_{\text{LEC}} \underbrace{U \partial_\delta U^+}_{\text{LEC}} = x_1 \\ - \partial_\beta U^+ U - U^+ \partial_\gamma U$$

etc.

$$\Rightarrow \mathcal{L}_{\text{break}}^{\text{ChPT}} = \frac{k}{2} C_{\alpha\beta\gamma\delta} \left[ \text{tr } \partial_\alpha U \partial_\beta U^+ \partial_\gamma U \partial_\delta U^+ + \text{tr } \partial_\alpha U^+ \partial_\beta U \partial_\gamma U^+ \partial_\delta U \right]$$

Set  $C_{\alpha\beta\gamma\delta}$  to physical value

$$\hookrightarrow \mathcal{L}_{\text{break}}^{\text{ChPT}} = \frac{k}{2} \cdot c \alpha^2 \left( \text{tr } \partial_\mu U \partial_\mu U^+ \partial_\nu U \partial_\nu U^+ + \text{tr } \partial_\mu U^+ \partial_\mu U \partial_\nu U^+ \partial_\nu U \right)$$

same terms

$$\boxed{\mathcal{L}_{\text{break}} = k \cdot c \cdot \alpha^2 \text{tr} (\partial_\mu U \partial_\mu U^+ \partial_\nu U \partial_\nu U^+)} \quad (1)$$

$$\text{Expansion in pion fields: } \propto (\partial_\mu \pi^\alpha \partial_\mu \pi^\alpha) (\partial_\nu \pi^\beta \partial_\nu \pi^\beta)$$

Power counting: If  $a \sim m_q \sim p^2 \Rightarrow$  (1) is of  $O(p^8)$

Very small!